



Fig. 10. Output voltage of the cavity. (a) Time response, horizontal scale 10 ns/cm, vertical scale 200 mV/cm with 60-dB attenuation in the line; (b) Spectral response, horizontal scale 100 MHz/cm (0-1 GHz coverage), vertical scale linear.

frequency of 2.1 GHz and 27 kW into  $50\ \Omega$  at a frequency of 1.5 GHz. The rise times were on the order of 200 ps. Using the available data [7] for normalized spark-gap resistance it was concluded that operation of the device at *X* band is feasible.

#### IV. CONCLUSIONS

For proper operation of this device at frequencies less than 3 GHz the spark resistance must fall to a value less than the characteristic impedance of the line in a time less than  $T$  where  $f_0 = 1/2T$  is the required frequency. The faster the rise time the more efficient is the operation of this device.

If a partial short is used, for minimal frequency shift due to coupling, the largest possible ratio of  $Z_0/Z_1$  (where

$Z_0$  is the characteristic impedance of the line and  $Z_1$  is the characteristic impedance of the output transmission line) should be used, and the inductance  $L_1$  should be adjusted for the required coupling. At low frequencies when the cross-sectional dimensions are limited it may not be possible to make  $L_1$  large enough to give the required coupling. In this case it will be necessary to decrease  $Z_0$  in order to increase the coupling. The decrease of  $Z_0/Z_1$  will result in a decrease of the resonant frequency. At frequencies higher than 3 GHz the spark-gap capacitance may be used to give a limited bandwidth output even though the normalized spark resistance,  $R$  is less than unity.

#### ACKNOWLEDGMENT

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#### Correction to "Behavior of the Magnetostatic Wave in a Periodically Corrugated YIG Slab"

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In the above short paper,<sup>1</sup> in (21) the minus sign was inadvertently left out. Equation (21) should read as follows:

$$-B_{z0} \cos \theta + B_{x0} \sin \theta = -B_z \cos \theta + B_x \sin \theta.$$

Consequently, the following correction should be noted.

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<sup>1</sup>M. Tsutsumi, Y. Sakaguchi, and N. Kumagai, *IEEE Trans. Microwave Theory Tech.*, vol. MTT-25, pp. 224-228, Mar. 1977.

1) In the second and seventh lines of (23) substitute  $+(\pi/d)\Delta$  and  $+(K_z + (2\pi/d)m)$  for  $-(\pi/d)\Delta$  and  $-(K_z + (2\pi/d)m)$ , respectively;

2) In (24) and (27), the factor  $[1 - ((\pi/d)\Delta)^2]$  and  $[\mu_1 - ((\pi/d)\Delta)^2]$  should have read  $[1 + ((\pi/d)\Delta)^2]$  and  $[\mu_1 + ((\pi/d)\Delta)^2]$ , respectively.

Since the numerical computation was based on the assumption that  $1 \gg ((\pi/d)\Delta)^2$  and  $\mu_1 \gg ((\pi/d)\Delta)^2$ , the results and conclusions reported in the paper are unaffected by this revision.